

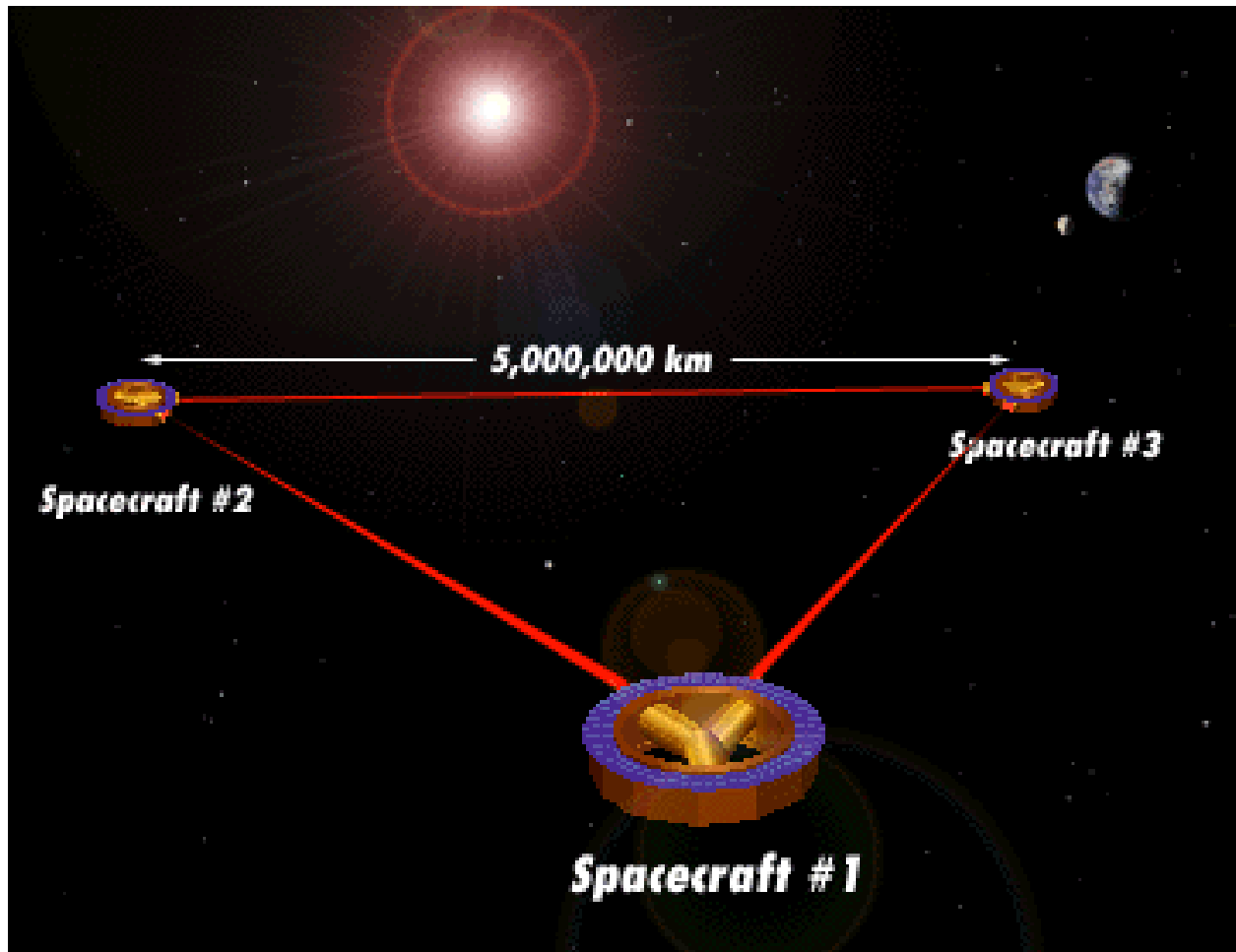
LISA: THE LASER INTERFEROMETER SPACE ANTENNA

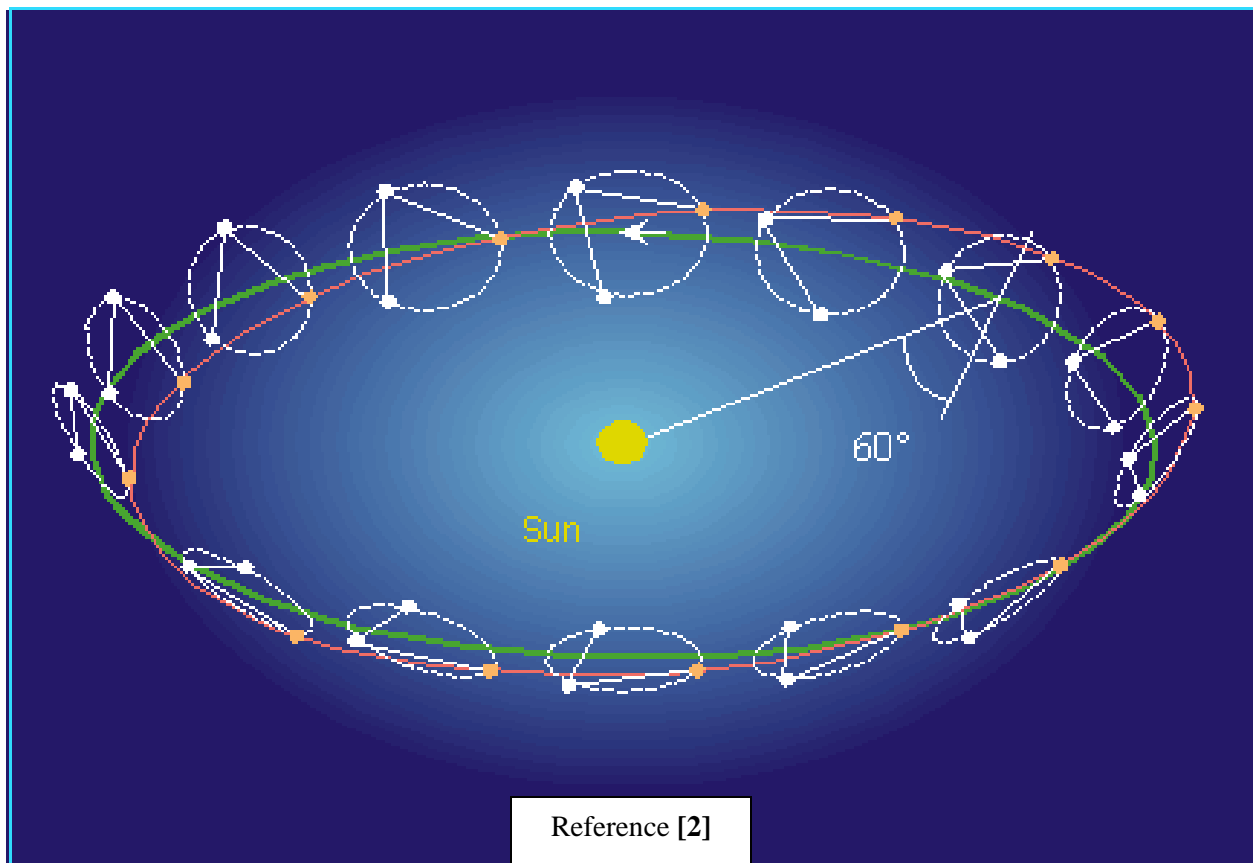
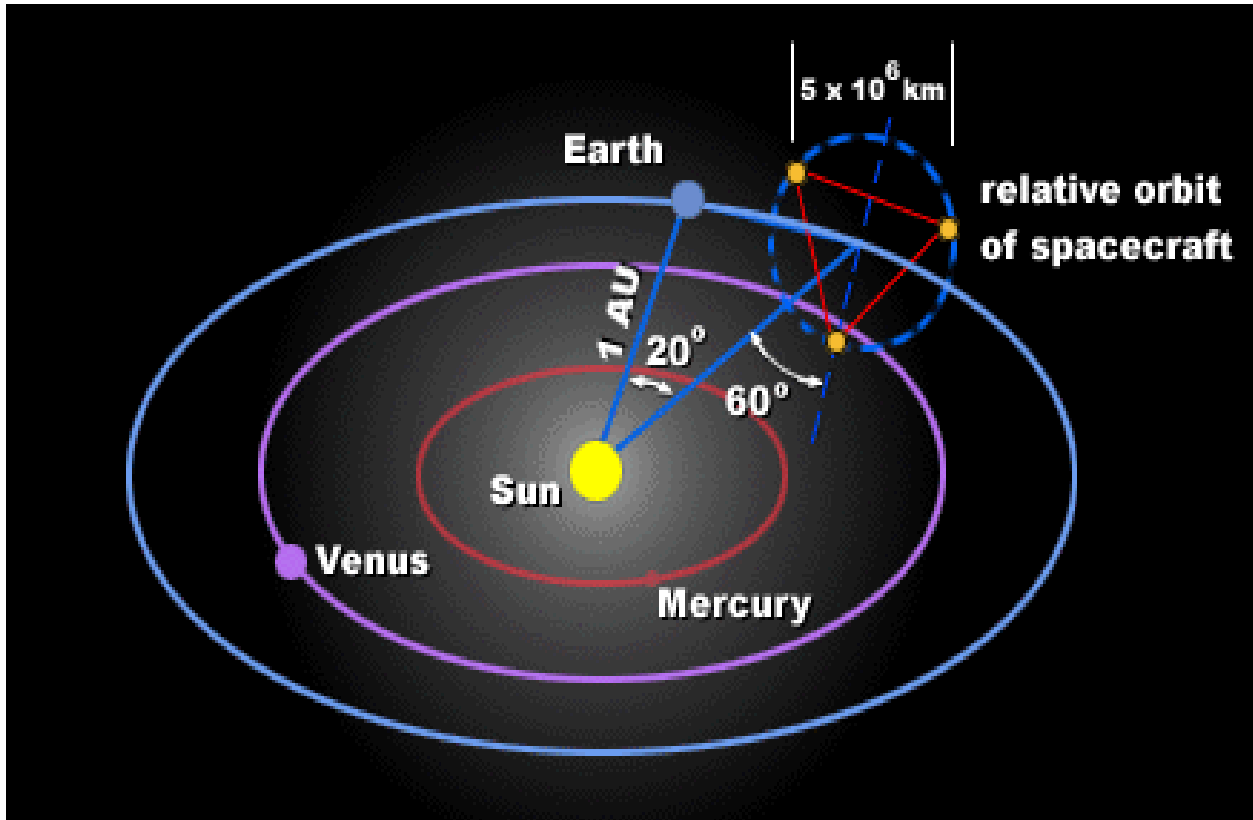
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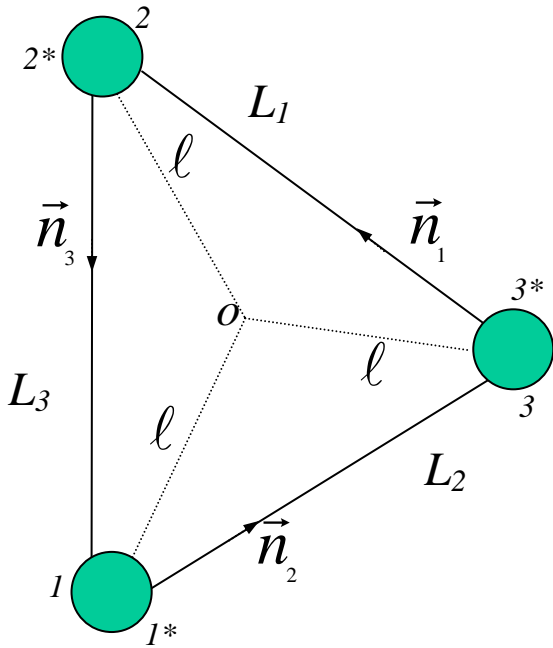
Abstract

The Laser Interferometer Space Antenna (LISA) is a deep-space mission, jointly proposed to the National Aeronautics and Space Administration (NASA), for detecting and studying gravitational radiation in the millihertz frequency band [1]. An overview of this new, exciting, and technologically challenging mission is presented, giving special emphasis to its frequency and timing requirements.



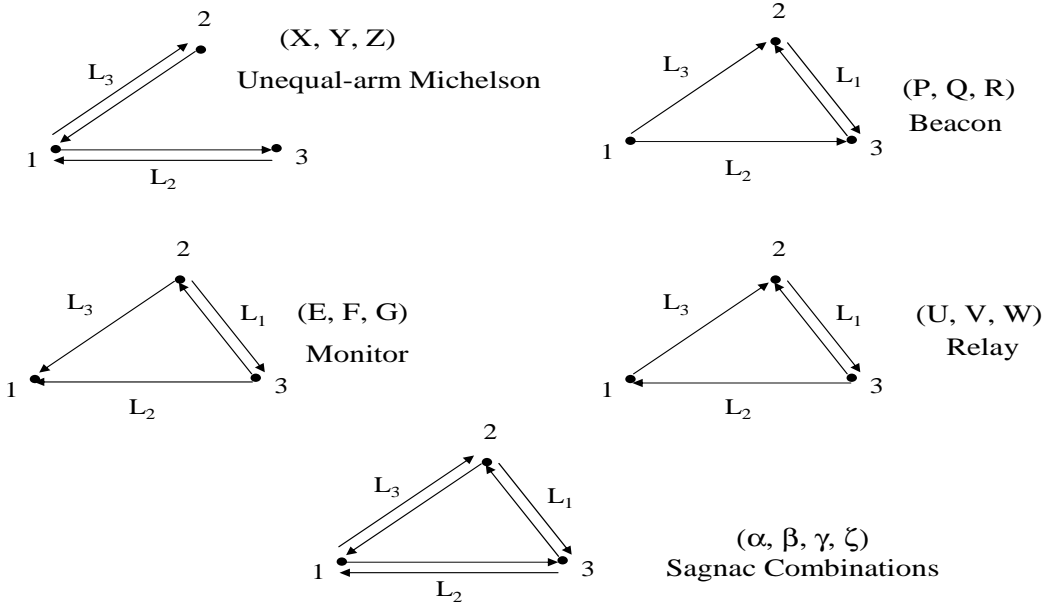


TIME-DELAY INTERFEROMETRY

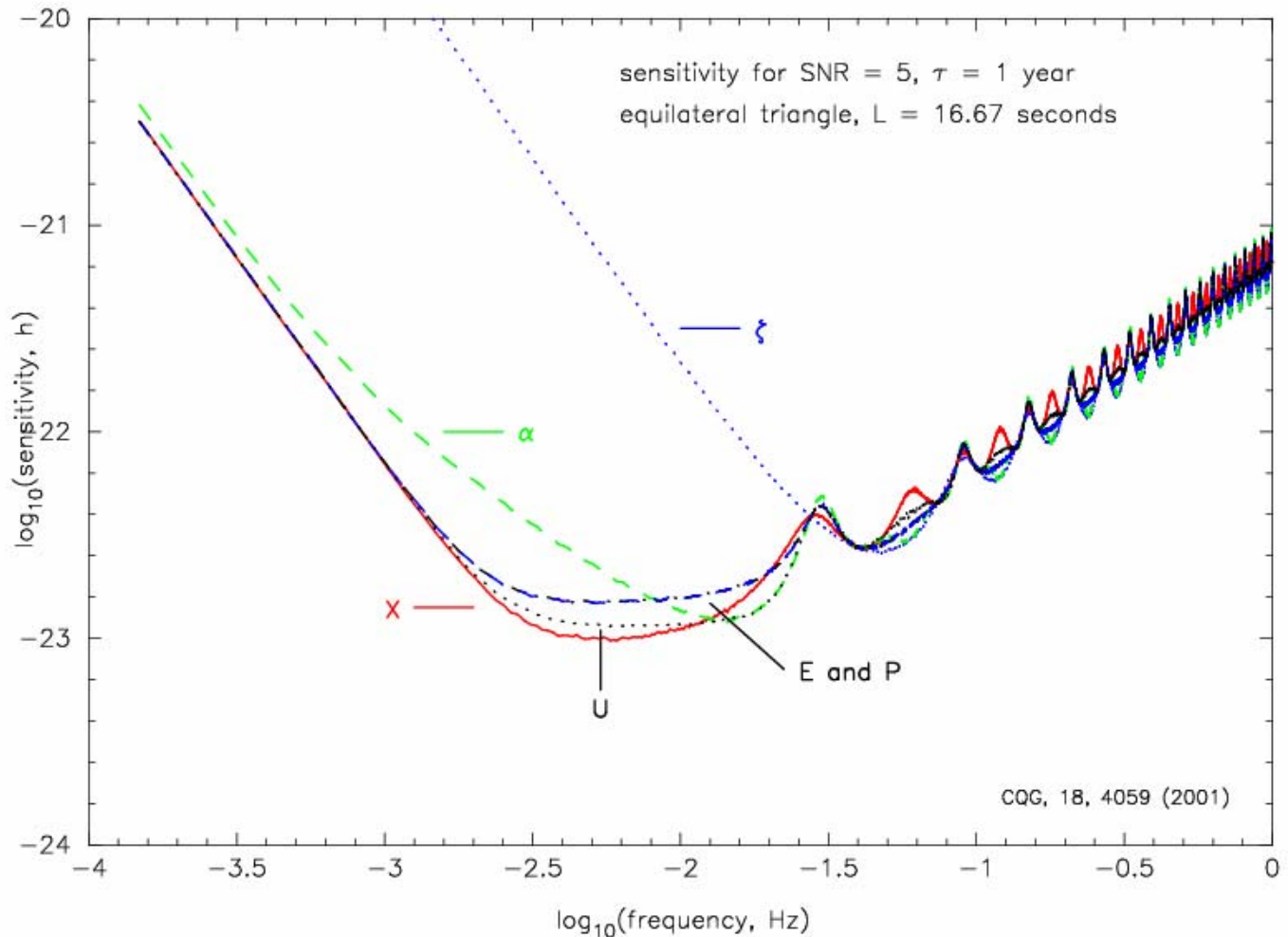


It is best to think of LISA as a closed array of six one-way delay lines between the test masses. This approach allows one to construct several interferometric data combinations. It offers advantages in hardware design, in robustness to failures of single links, and in redundancy of data [3-5].

INTERFEROMETRIC DATA COMBINATIONS



SENSITIVITIES

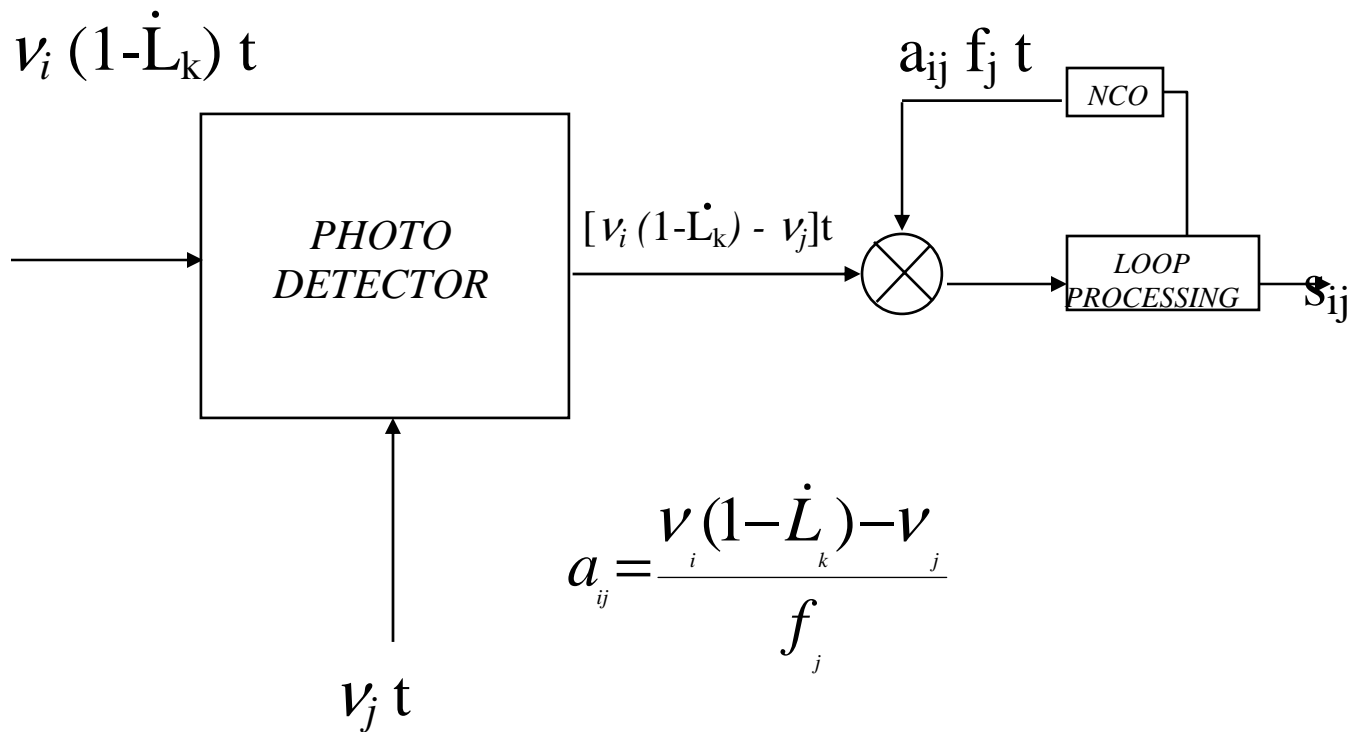


MOVING SPACECRAFT ARRAYS AND CLOCKS SYNCHRONIZATION

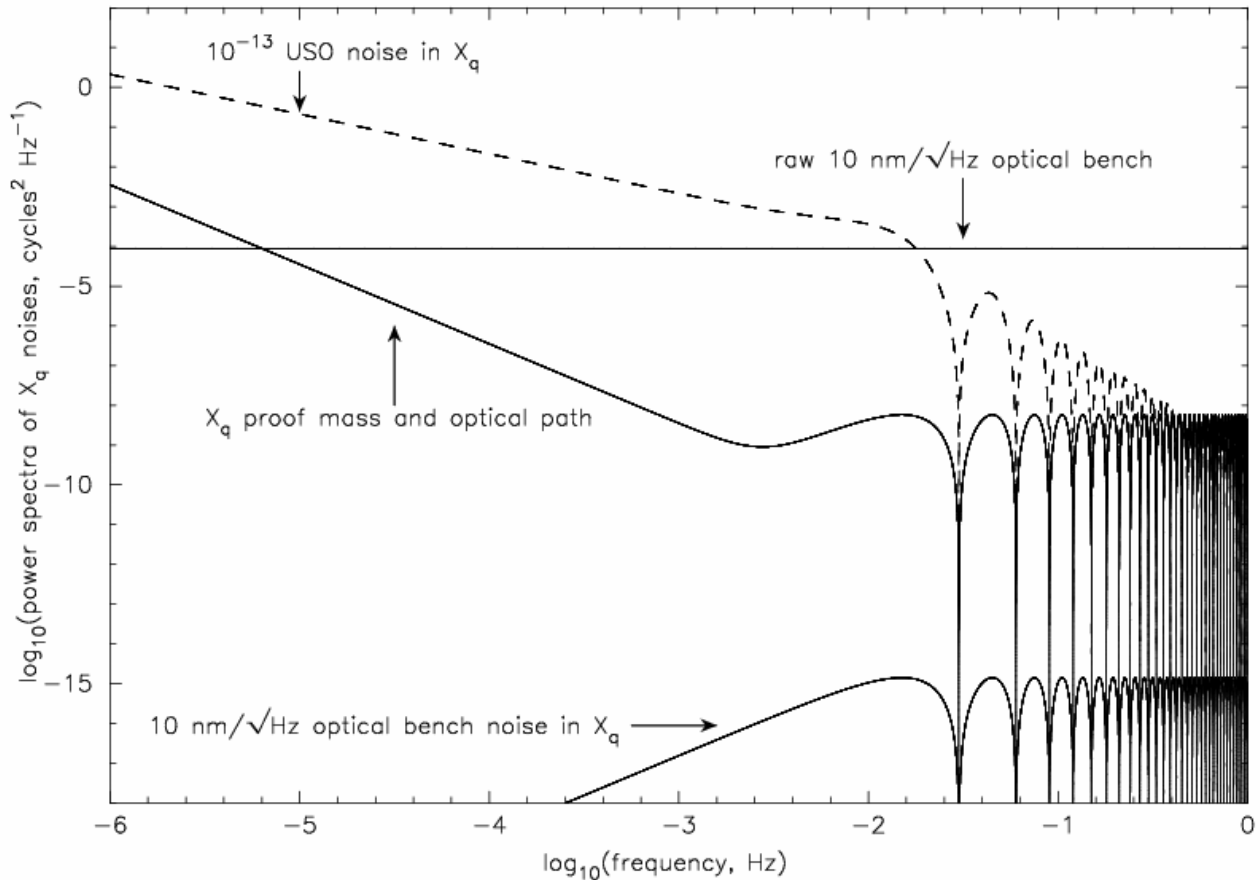
The clocks onboard the LISA S/Cs ought to be synchronized to each other with an accuracy of 40 ns or better in order to suppress the lasers frequency fluctuations below the level of the remaining noise sources. In a rotating reference frame, the Sagnac effect prevents the implementation of the Einstein's Synchronization Procedure, i.e. synchronization by transmission of electromagnetic signals (GPS is a good example of this problem!) To account for the Sagnac effect, one introduces a *hypothetical* inertial reference frame, and time in this frame is the one adopted by the spacecraft clocks! In other words, the onboard receivers have to convert time information received from Earth to time in this inertial reference frame (SSB) [6-7].

In the SSB frame, the differences between back-forth delay times are very much larger than has been previously recognized. The reason is in the aberration due to motion and changes of orientation in the SSB frame. With a velocity $V=30$ km/s, the light-transit times of light signals in opposing directions (L_i and L'_i) will differ by as much as $2VL$ (a few thousands km). They will also change in time due to rotation (0.1 m/s); this however is significantly smaller than the spacecraft relative velocity (10 m/s). The arm lengths need to be known with an accuracy of 100 m or better for suppressing the lasers frequency fluctuations below the level of the remaining noises [8-9].

The laser frequencies *are* different, and the spacecraft *are not* stationary. Both frequency offsets between lasers, and Doppler drifts, now bring in noise from the onboard oscillators (USO=ultra stable oscillators) used in the down-conversion of phototube fringe rates [10].



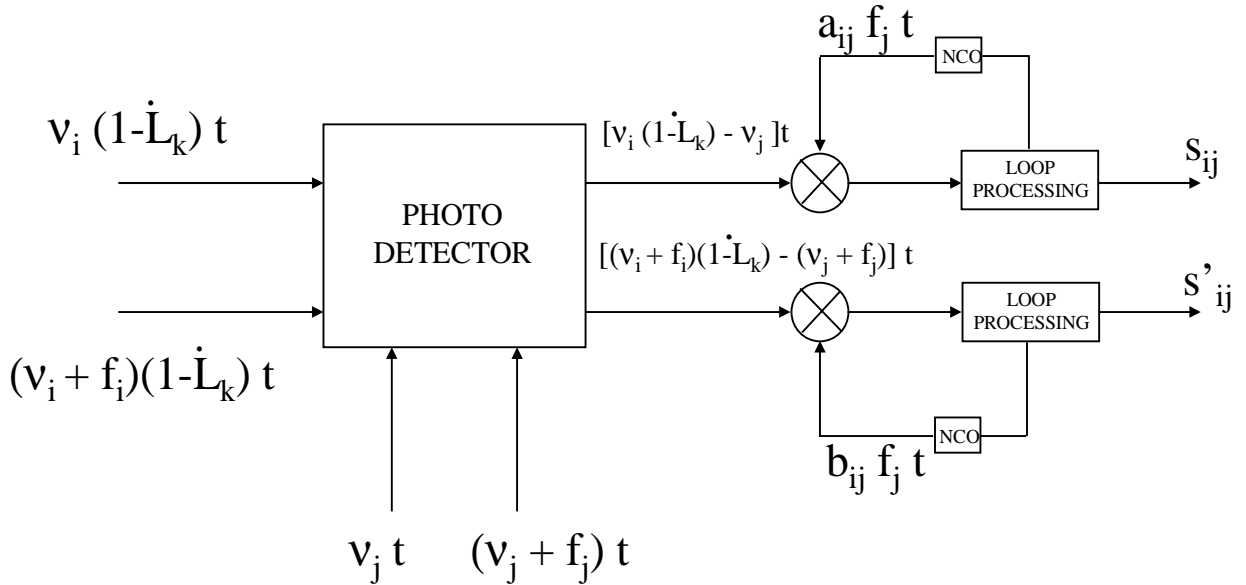
INTERFEROMETRIC COMBINATIONS



USO NOISE CALIBRATION

In the scheme first proposed by Hellings et al. [11], in addition to the six main laser signals of frequencies n_i , a second laser signal is superimposed on each beam by either modulating it at the frequency f_i of its USO (two main side-bands), or by combining each beam with a coherent second signal at $n_i + f_i$.

The transmitted second signals are heterodyned against the local second signal, and independently down converted with coefficients, say b_{ij} (different from the a_{ij} introduced earlier!) to give six additional data records, s'_{ij} . By introducing the observables $r_{ij} = (s_{ij} - s'_{ij})/f_k$, USO phase noise can be measured and calibrated-out from the interferometric combinations.



$$a_{ij} = \frac{v_i(1-\dot{L}_k) - v_j}{f_j} \quad ; \quad b_{ij} = \frac{(v_i + f_i)(1-\dot{L}_k) - (v_j + f_j)}{f_j}$$

CONCLUSIONS

- T.D.I. provides a robust method for canceling the leading noise source – laser phase fluctuations – in an interferometer with unequal, time-variable arms.
- The onboard clocks have to be synchronized with an accuracy of 40 ns or better in order to synthesize interferometric measurements.
- The heterodyne phase measurements require an onboard USO on each spacecraft, and a highly precise and accurate Phase Meter.

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QUESTIONS AND ANSWERS

TOM CLARK (Syntonics): The solar system has many, many bodies in it, each one of which has some natural periodicity in its gravity signature. I would think that the one over F noise in that combined set of thousands of bodies would have quite a bit of power at these low frequencies like 200 seconds that you were describing for the gravity signature.

MASSIMO TINTO: You are talking in terms of as a gravitational receiver or as a perturbation to this system?

CLARK: Just as gravitational noise.

TINTO: Something that I didn't mention actually, but in our galaxy there are hundreds of millions of binary systems, white dwarf binaries, which, actually, we allow to appear incoherently to the measurements, in at least a weaker form. In balancing, in a sense, it will provide the gravitational wave confusion noise. They will sit above the sensitivity level that we estimated.

So, from one point of view, it can be regarded as a beautiful signal so you will detect it and you will see it. If you are able to discriminate it against the noise, you will have a detection. But at the same time, if you want to observe other kinds of signals, then you have to account for it and perhaps try to remove it.

